# Apodization of Coupled Cavity Array for Waveguide Quantum Electrodynamics: Design & Simulation

# – Hybrid Quantum Circuit Laboratory (HQC) –

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submitted by

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## <span id="page-2-0"></span>1 Introduction and Motivations

In the present work, we aim to optimize the design of a coupled resonator waveguide QED to obtain flat transmission on a defined frequency range through apodization of the transmission peaks. In order to reach apodization we slightly modify the physical parameters of the resonators at both ends of the waveguide device so that the reflections cancel out the Fabry–Perot resonances at the resonant frequencies. This leads to a broadening of the characteristic peaks usually observed in transmission for coupled cavity arrays. Under optimal conditions, investigated by Chak et al. [\[1\]](#page-15-0), we can exploit these broadenings to obtain flat transmission on a frequency band whose width depends on the number of resonators in the waveguide and on their coupling strength.

Apodized coupled resonator waveguides have been used for filtering applications as they operate as band-pass filters [\[2\]](#page-15-1). However, in recent years these devices have captivated more interest because of their possible use in slow light applications [\[3\]](#page-15-2). Slow light devices are based on the idea that the dispersion relation in periodic structures flattens at the transmission band edges. Since the group velocity is defined by the derivative of the dispersion relation, it is therefore small in regions where the dispersion relation goes flat, meaning that around the transmission band edges we can achieve both large transmission and low group velocity. These devices can be used to study non-Markovian dynamics i.e. dynamics of physical systems with memory of the past interactions [\[4\]](#page-15-3).

Superconducting coplanar waveguide resonators have been also used to achieve strong coupling regime, paving the way for a multitude of new investigations [\[5,](#page-15-4) [6\]](#page-15-5). In addition, they represent an ideal platform to study the quantum dynamics of systems interacting with a continuum of EM modes (multimode coupling).

### <span id="page-3-0"></span>2 Materials and Methods

In the following section, I am going to provide a detailed account of the procedure that was followed in completing the project and the tools that were adopted. After a concise review of the theoretical background in [subsection 2.1,](#page-3-1) I will present the software tools used for simulations [\(subsection 2.2\)](#page-5-0) and, finally, the workflow that guided the entire project [\(subsection 2.3\)](#page-5-1).

#### <span id="page-3-1"></span>2.1 Apodization in arrays of lumped-element resonators

We consider an array of  $n = 8$  resonators, all featuring the same parameters, except for the first two and the last two.



**Figure 1:** Coupled resonators waveguide schematic.  $\omega_0$  and  $\omega'_0$  are the resonators resonant frequencies while J and J' are the coupling strengths.  $\gamma$  is the energy decay rate of the edge sites

We describe the resonator array using the tight-binding formalism, assuming periodic boundary conditions and considering only nearest-neighbor coupling. The corresponding n-site Hamiltonian is the following:

$$
H = \begin{pmatrix} \omega'_0 + i\gamma & J' & 0 & & & \\ {\bf J'} & \omega_0 & {\bf J} & & & \\ {\bf 0} & {\bf J} & \omega_0 & & & \\ & & \ddots & & & \\ & & & \omega_0 & {\bf J} & {\bf 0} \\ & & & & {\bf J} & \omega_0 & {\bf J'} \\ & & & & {\bf 0} & {\bf J'} & \omega'_0 + i\gamma \end{pmatrix}
$$

Where  $\omega_0'$  is the first (and last) resonator resonant frequency, J is the inter-site coupling, J' is the inter-site coupling between the first two (and the last two) resonators,  $\gamma$  is the energy decay rate of the edge sites due to the coupling to the measurement setup. We realize this Hamiltonian with a lumped-element resonator array coupled to transmission lines (with  $Z_0$  feedline input impedance) whose equivalent circuit is showed in [Figure 4.](#page-7-1)

The lumped-element resonator array parameters are related to the Hamiltonian parameters by:

 $J = \frac{C_c \omega_0}{2C_E}$  $\frac{C_c\omega_0}{2C_{\Sigma}}$ : Inter-site coupling with total capacitance  $C_{\Sigma} = C + 2C_c$  and resonant frequency  $\omega_0 = \frac{1}{\sqrt{L}d}$  $LC_{\Sigma}$ 

 $J' = \frac{C_{c1}\omega'_0}{2C_{\Sigma'}}$ : First two (and last two) resonators coupling with total capacitance  $C_{\Sigma'} = C_k +$  $2C_2 + C_{c1}$  and resonant frequency  $\omega'_0 = \frac{1}{\sqrt{L}}$  $LC_{\Sigma'}$ 

 $\gamma = \frac{C_k^2 Z_0 \omega_0}{C^2 Z}$  $\frac{C_E^2 Z_0 \omega_0}{C_\Sigma^2 Z_r}$ : Energy decay rate of the edge sites, with  $Z_r = \sqrt{L/C_\Sigma}$ 

In order to reach optimal apodization and obtain a flat passband in transmission, the conditions proposed by Chak and coworkers are the following [\[1\]](#page-15-0):

$$
\omega'_0 = \omega_0
$$
  

$$
J' = \sqrt{2}J
$$
  

$$
\gamma/2 = 2J\sin(\psi)
$$

where  $\psi = \pi/2$  (best matching at the middle of the band) or  $\psi = \pi/4$  (best matching at the band edges). Considering the limit  $C > C_k >> C_{c1}$ ,  $Cc$ , we can meet these conditions with the following capacitances values:

$$
C_{c1} = \sqrt{2}C_c
$$
  
\n
$$
C_k = \sqrt{\frac{2C_c sin(\psi)}{Z_{0}\omega_0}}
$$
  
\n
$$
C_1 = C - (C_{c1} - C_c)
$$
  
\n
$$
C_2 = C - (C_{c1} + C_k - 2C_c)
$$

Where the first two conditions optimize the capacitance values to have the best result, while the last two ensure that  $\omega'_0 = \omega_0$ .



Figure 2: Equivalent circuit representing the resonators waveguide on lumped element description. Note that we consider only nearest-neighbor coupling, secondary coupling has been neglected.

In [section 3](#page-13-0) LTspice simulations have been used to prove that the values derived through these equations give indeed the best result: even with a small number of resonators, flat transmission is observed on a well-defined frequency range proportional to the number of resonators.

#### <span id="page-5-0"></span>2.2 Softwares and tools

In the current subsection, I am going to briefly present the software used for simulations and the Python library used to design the resonators waveguide.

LTspice and Cadence Microwave Office: circuit design softwares. Highperformance SPICE simulation software, schematic capture and waveform viewer with enhancements and models for easing the simulation of analog circuits.

Sonnet Software: EDA software solutions. Note that, despite Sonnet and LTspice have been used to calculate the transmission spectrum, LTspice simulations are based on a circuit diagram that models the device while Sonnet simulates by numerically calculating Maxwell equations point by point on a 2-D representation of the real design. for high-frequency RF/MW electromagnetic analysis.

ANSYS: 3D electromagnetic (EM) simulation software for designing and simulating high-frequency electronic devices.

gdspy library: Python module that allows the creation of GDSII stream files for micro and nanofabrication.

#### <span id="page-5-1"></span>2.3 Workflow

The following subsection is dedicated to the presentation of the project workflow: I will provide an overview of the resonator waveguide design process, starting from the equations and finishing with the final waveguide design that will be fabricated.

#### <span id="page-5-2"></span>2.3.1 Optimal capacitance values determination

The first step consisted in exploiting the equations, derived in [subsection 2.1,](#page-3-1) in order to find the best capacitance values for the resonator waveguide to obtain a flat band in transmission. Since we have more parameters than equations: 6 parameters to determine

$$
C, C_c, C_{c1}, C_k, C_1, C_2
$$

and only 4 equations,

$$
C_{c1} = \sqrt{2}C_c
$$
  
\n
$$
C_k = \sqrt{\frac{2C_c sin(\psi)}{Z_0 w_0}}
$$
  
\n
$$
C_1 = C - (C_{c1} - C_c)
$$
  
\n
$$
C_2 = C - (C_{c1} + C_k - 2C_c)
$$

we are free to assign arbitrary values to two of these parameters and successively derive the others. This comes really in handy when considering that, in circuit quantum electrodynamics, we have a narrow range of possible values for these capacitances due to the physical limits of the devices. We can therefore exploit these degrees of freedom to choose some convenient values that are easy to realize.

Note that, in the present analysis, we are not considering  $L_0$  and  $Z_0$  since these parameters are fixed.  $Z_0$  is the input feedline impedance for which we have only two possible values: 50Ω, and 500Ω which can be reached by tapering. On the other hand, we take  $L_0$  as a fixed parameter  $(L_0 = 38.8nH)$  because we want a large inductance since it is needed for metamaterial waveguides  $[7, 8, 9]$  $[7, 8, 9]$  $[7, 8, 9]$  $[7, 8, 9]$  $[7, 8, 9]$ , paving the way for future applications of the devices designed in this work. When increasing  $L_0$  we induce a redshift in the resonators' resonant frequency (defined as  $w_0 = \frac{1}{\sqrt{L}}$  $\frac{1}{LC_{\Sigma}}$ ), therefore we cannot use inductance values larger than  $\approx 40 \text{ nH}$ , otherwise the resonant peaks in transmission would appear in a range of the frequency spectrum we are not able to probe with the available lab instruments.

I started with  $Z_0 = 50\Omega$  and chose as starting point the following values:

$$
C_c = 0.43 \text{ } fF \text{ } , \text{ } C = 25 \text{ } fF
$$

where  $C_c$  represent the coupling capacitance between two neighboring resonators in the waveguide and C is the coupling capacitance between a single resonator and ground since these values have been already obtained for other devices in the past, therefore we knew they were achievable in practice.

Using Chak's equations we successively derived the values for the other parameters:

	0.43 fF   25 fF   19.63 fF   0.61 fF   24.82 fF   5.62 fF		

**Table 1:** Values derived through Chak's equations starting from  $C_c = 0.43$  fF,  $C =$ 25 fF

#### <span id="page-6-0"></span>2.3.2 LTspice/Cadence simulations and design code

Once a proper choice of the starting parameters has been made, I simulated the transmission spectrum of the resonator waveguide using LTspice and Cadence Microwave Office. Simulations and parametric analysis have been used to confirm that the derived parameter values were indeed the ones giving the best results. As it is shown in [Figure 4,](#page-7-1) the closer the parameters are to the derived ones, the flatter is the band we observe in transmission.

After the simulations confirmed this starting set of values, I designed the first prototype of the device. In order to optimize and speed up the designed process I coded a script in Python, using the *qdspy* library, that prints out a fully customizable gds file of the resonators waveguide (the Python code is reported in [Appendix A\)](#page-17-0)

<span id="page-7-2"></span>

<span id="page-7-1"></span>Figure 3: Transmission vs frequency for a standard coupled resonators waveguide not apodized (the parameters are set as:  $C_c = C_{c1} = C_k$  and  $C_1 = C_2 = C$ ).



Figure 4: Transmission spectrum for the apodized coupled resonators waveguide. The 4 traces in different colors represent 4 different sets of values for the parameters  $C_k$ and  $C_{c1}$ . It is clear that the best result is achieved when using the values derived from Chak's equations: for  $C_k = 19.63$  fF and  $C_{c1} = 0.61$  fF a flat band is observed in transmission.

#### <span id="page-7-0"></span>2.3.3 Design optimization with  $Z_0 = 50 \Omega$

The gds file has been then used to run a simulation on ANSYS to numerically calculate the capacitance matrix of this first design. The capacitance matrix provides, for each element of the waveguide (single resonators, ground, feedline capacitors), the coupling capacitances with respect to the other elements. The capacitance matrix values have

been then compared to the desired capacitance values derived from the equations. This first comparison allowed me to understand how far the starting design was from the optimal one I was looking for.

<span id="page-8-0"></span>

Figure 5: First waveguide design: the ground piece between the first (last) and second (last second) resonators has been removed to reduce the capacitance to ground for the first resonator  $(C_2)$ . Note that, for ANSYS simulations, the resonators must be separated from the ground (as you can see in this picture) so that they can be simulated as individual elements.

			$\smile_{c1}$	
Desired cap values	$0.43$ fF	25 fF   19.63 fF   0.61 fF   24.82 fF   5.62 fF		
First design cap values $ \approx 1.5$ fF $ \approx 17$ fF $ \approx 5.5$ fF $ \approx 2.5$ fF $ \approx 15$ fF $ \approx 14.5$ fF				

Table 2: Comparison between ANSYS simulation results and optimal capacitance values derived from Chak's equation

The next phase has been characterized by a trial and error approach aimed at understanding how different modifications in the physical structure of the device affected the capacitance values. To this purpose, I have run several simulations on ANSYS changing each time a different parameter in the waveguide: for example capacitors thickness, capacitors length, distances between resonators and ground, and so on.

This analysis showed that, as expected, we can modify the coupling capacitance between two elements in the waveguide by:

- Increasing or decreasing the distance between the two elements. However, we cannot have a too large separation since we would have problems with the lithography process (used to fabricate the waveguide): first of all, it would imply longer time and higher costs for the process, and it would also increase the probability of errors in the making.
- Modifying the thickness of the two elements. However, we cannot decrease it too much or an element that is supposed to be a capacitor will start behaving as an inductor in practice.
- Modifying the exposed surface between the two elements (see example in [Fig](#page-9-0)[ure 6\)](#page-9-0).

<span id="page-9-0"></span>

**Figure 6:** Impact of the feedline capacitor length on  $C_k$  value.  $C_k$  represents the coupling capacitance between the first resonator and the feedline capacitor. Note that there are other waveguide features that strongly affect the value of  $C_k$ : the distance between the feedline capacitor and the resonator capacitor, and the thickness of the capacitors.

After this initial study, I started modifying the waveguide design in order to achieve the optimal capacitance values. This step by step process gradually converged towards the optimal waveguide design: at each step, the design was modified and simulated with ANSYS, then the simulations' results, compared with the desired capacitance values, provided feedback that suggested how to further modify the design.

<span id="page-9-1"></span>

Figure 7: Sonnet simulation of the first waveguide design. The transmission spectrum shows several well defined peaks meaning that the waveguide is not apodized.

As shown [Figure 5,](#page-8-0) the ground between the first and second resonators has been removed in order to decrease the capacitance to ground for the first resonator with respect to the others (indeed for the optimal capacitance values we have  $C_2 \ll C$ ) and increase the capacitive coupling between the first two resonators (since  $C_{c1} > C_{c}$ ).

During this design optimization process, together with ANSYS simulations, Sonnet simulations [\(Figure 7\)](#page-9-1) were used to monitor the evolution of the transmission spectrum.

Thanks to this process I managed to design a waveguide that met the optimal parameters for  $C, C_c, C_c, C_k, C_1$ . However, the simulations showed that, for this kind of device, obtaining  $C_2 = 5.62$  fF was not feasible (where  $C_2$  is the coupling capacitance between the first/last resonators and ground). Even pushing the design to the limits

the  $C_2$  value calculated by ANSYS did not get smaller than  $\approx 10-11$  fF [\(Figure 8\)](#page-10-0). I could not even lower the values of  $C_c$  and  $C_{c1}$  (where  $C_{c1} = \sqrt{2C_c}$ ) to increase the optimal value of  $C_2$  (where  $C_2 = C - (C_{c1} + C_k - 2C_c)$ ) because they both have a very narrow range of achievable values ( $\approx 0.1 - 3$  fF) and they were already quite small. If  $C_c$  and  $C_{c1}$  are too small, the waveguide does not work properly since these two values are proportional to the capacitive coupling between the resonators  $J \propto C_c$  and  $J' \propto C_{c1}$ .

<span id="page-10-0"></span>

**Figure 8:** Attempts to decrease  $C_2$  (first and last resonators' capacitance to ground) in order to achieve the optimal value. In designs  $(a)$  and  $(b)$  the separation between the first resonator and the ground is increased, in  $(c)$  and  $(d)$  the separation is increased and the first resonator capacitor thickness is reduced. None of the proposed designs reach the optimal value for  $C_2$  and it is not possible to increase further the distance from the ground since this would give problems during fabrication with the lithography process.

In order to overcome this problem, the solution was finding a new optimal capacitance values set. Therefore I exploited again the equations presented in [subsection 2.1,](#page-3-1) imposing  $C_2 = 13$  fF this time. Note that  $C_2$  has to be kept as small as possible because it is directly related to C value: following Chak's equations, the larger  $C_2$ the larger C is gonna be. Already for  $C_2 \approx 15{\text -}20$  fF, C reaches values that are not achievable in practice for this kind of devices. Imposing  $C_2 = 13$  fF allows to have achievable values for both  $C_2$  and  $C$ .

1.00 fF   48.00 fF   35.29 fF   1.41 fF   47.58 fF   13.30 fF		

<span id="page-10-1"></span>**Table 3:** New set of optimal values derived starting from  $C_2 = 13$  fF and  $C_c = 1$  fF.

Starting from  $C_2 = 13$  fF and  $C_c = 1$  fF, which ensures a good coupling between the resonators, I derived the new set of optimal capacitance values reported in [Table 3.](#page-10-1)

This new set contains all achievable capacitance values. However, in this case we have the resonators' resonant frequencies approaching  $4GHz$ :

$$
f_0 = \frac{1}{\sqrt{LC_{\Sigma}} 2\pi} \approx 3.6 \; GHz
$$

This is a problem for us because we cannot probe this range of frequencies with the instruments available in our lab.

<span id="page-11-0"></span>

**Figure 9:** ANSYS Capacitance matrix of the optimal design for the case  $Z_0 = 50 \Omega$ 

The only remaining parameter that can be tuned is  $L_0$ , even though reducing  $L_0$  means giving up on high inductance to ground for our device. We have that:

$$
L=L_\square\frac{l}{w}
$$

Where  $L_{\Box}$  is the inductance per square, l and w are the length and the width of the inductor wire. By shortening the inductor wire by almost half of its original length, I reduced  $L_0$  from 38.8 nH to 20 nH, taking the resonator's resonant frequencies to  $\approx 5$ GHz.

With this new set of parameters, adopting the same optimization procedure described before, I obtained the design shown in [Figure 9.](#page-11-0) Note that the thickness of the resonators' capacitors has been increased and the distance between the resonators and ground has been strongly decreased to achieve large C values. On the other hand, reaching a large  $C_k$  has been easy since the interdigitated feed line capacitor allows good control over this parameter.

Note also that, despite the capacitance values calculated through ANSYS simulation (shown in [Figure 9\)](#page-11-0) are not close to the optimal values (shown in [Table 3\)](#page-10-1), this design is the one giving the best results for the case  $Z_0 = 50 \Omega$ .

#### <span id="page-12-0"></span>2.3.4 Design optimization with  $Z_0 = 500 \Omega$

It is possible to achieve an input impedance of  $Z_0 = 500 \Omega$  by tapering. This second case turned out to be much easier since a larger value of  $Z_0$  allows to reach apodization with lower capacitance values with respect to the previous case, which is easier to achieve, while maintaining at the same time a large  $L_0$  ( $L_0 = 38.8$  nH).

<span id="page-12-1"></span>Following the same process used for the  $Z_0 = 50 \Omega$  case, I derived the set of optimal values reported in [Table 4.](#page-12-1)

L 00 fF		22.70 fF   9.35 fF   1.41 fF   22.20 fF   13.98 fF	

Table 4: New set of optimal values for the case  $Z_0 = 500 \Omega$ 

<span id="page-12-2"></span>

**Figure 10:** ANSYS Capacitance matrix of the optimal design for the case  $Z_0 = 500 \Omega$ 

Again, we obtained the design reported in [Figure 10](#page-12-2) for this new case by gradual design optimization following the steps described previously.

### <span id="page-13-0"></span>3 Results and Conclusions

In the following section, the final results and the conclusions are presented: I will show the two final designs, products of the optimization process presented in [subsection 2.3,](#page-5-1) and the relative transmission spectra.

#### <span id="page-13-1"></span>3.1 Results

<span id="page-13-2"></span>The optimal design for the case  $Z_0 = 50 \Omega$  is reported in [Figure 11:](#page-13-2)



**Figure 11:** Optimal design (a) and transmission spectrum (b) for the case  $Z_0 = 50 \Omega$ 

For the second case,  $Z_0 = 500 \Omega$ , the results are shown in [Figure 12:](#page-13-3)

<span id="page-13-3"></span>

Figure 12: Optimal design (a) and transmission spectrum (b) for the case  $Z_0 = 500 \Omega$ 

In [Figure 11](#page-13-2) and [Figure 12](#page-13-3) we can observe that both the transmission spectra feature an asymmetric shape, a shoulder is visible on the right side of the bands. This asymmetry is inherited from the resonator waveguide transmission spectrum, which already shows this characteristic when not apodized (see [Figure 3\)](#page-7-2).

Another common characteristic is the presence of two residual peaks on the left side of the bands. We observe that these peaks are less pronounced in the second case  $(Z_0 = 500 \Omega)$ , where we also have a larger frequency range with flat transmission:  $\approx 0.3$  GHz, compared to  $\approx 0.15$  GHz for the  $Z_0 = 50 \Omega$  case. In order to mitigate these oscillations at the band edges we could use a larger number of resonators: with more resonators we can increase the extent of the band and obtain a flat transmission on a larger frequency range.

In [Figure 13](#page-14-1) are pictured the two different resonator designs for the two waveguides. In the first case  $(Z_0 = 50 \Omega)$  the capacitor has a wider and thicker bottom part in order to reach a larger value for the capacitance to ground  $C = 48$  fF, while  $C = 22.7$  fF in the case  $Z_0 = 500 \Omega$ . On the other hand, in order to compensate this larger capacitance value, the inductor wire is reduced by half in this case so that the resonators resonant frequencies (defined as  $f_0 = \frac{1}{2\pi\pi\sqrt{3}}$  $\frac{1}{2*\pi\sqrt{LC_{\Sigma}}}$ , where C is the dominant factor in  $C_{\Sigma} = C + 2C_{c}$ falls into the range  $4.5$  GHz -  $6\overline{\smash{\mathrm{GHz}}}$ .

<span id="page-14-1"></span>

Figure 13: Comparison between the two different resonators designs for the two waveguides

#### <span id="page-14-0"></span>3.2 Conclusions

Starting from the equations derived in [subsection 2.1,](#page-3-1) we successfully designed two coupled resonators waveguides, with different input impedances  $Z_0 = 50 \Omega$   $Z_0 =$ 500  $\Omega$ , optimized to achieve apodization. Both devices clearly show flat transmission respectively in the frequency ranges  $5.6 - 5.8$  GHz ( $Z_0 = 50 \Omega$  case) and  $5 - 5.3$ GHz ( $Z_0 = 500 \Omega$  case), according to Sonnet simulations. These frequency ranges are optimal to study waveguide-SC qubits multimode coupling since superconducting qubits usually have frequencies in the range  $4 - 8$  GHz [\[10,](#page-15-9) [11,](#page-15-10) [12\]](#page-16-0). The waveguide designs have been derived through the process explained in [subsection 2.3.](#page-5-1) In order to optimize and speed up the design process we developed a Python program that prints out a fully customizable gds file of the resonators waveguide (see [Appendix A\)](#page-17-0).

### References

- <span id="page-15-0"></span>[1] Philip Chak and J. E. Sipe. "Minimizing finite-size effects in artificial resonance tunneling structures". In: *Opt. Lett.*  $31.17$  (Sept. 2006), pp. 2568–2570. DOI: [10.](https://doi.org/10.1364/OL.31.002568) [1364/OL.31.002568](https://doi.org/10.1364/OL.31.002568). url: [http://www.osapublishing.org/ol/abstract.](http://www.osapublishing.org/ol/abstract.cfm?URI=ol-31-17-2568) [cfm?URI=ol-31-17-2568](http://www.osapublishing.org/ol/abstract.cfm?URI=ol-31-17-2568).
- <span id="page-15-1"></span>[2] J. Capmany et al. "Apodized coupled resonator waveguides". In: Optics express 15 (Sept. 2007), pp. 10196–206. doi: [10.1364/OE.15.010196](https://doi.org/10.1364/OE.15.010196).
- <span id="page-15-2"></span>[3] Simone Gasparinetti Simon Mathis Jean-Claude Besse. "Microwave Waveguides with Engineered Dispersion based on Arrays of Lumped-Element Resonators". In: ETH Zurich, semester project (Oct. 2017).
- <span id="page-15-3"></span>[4] Vinicius S. Ferreira et al. "Collapse and Revival of an Artificial Atom Coupled to a Structured Photonic Reservoir". In: Phys. Rev. X 11 (4 Dec. 2021), p. 041043. doi: [10.1103/PhysRevX.11.041043](https://doi.org/10.1103/PhysRevX.11.041043). url: [https://link.aps.org/doi/10.](https://link.aps.org/doi/10.1103/PhysRevX.11.041043) [1103/PhysRevX.11.041043](https://link.aps.org/doi/10.1103/PhysRevX.11.041043).
- <span id="page-15-4"></span>[5] Abdufarrukh A. Abdumalikov et al. "Vacuum Rabi splitting due to strong coupling of a flux qubit and a coplanar-waveguide resonator". In: Phys. Rev. B 78 (18 Nov. 2008), p. 180502. doi: [10.1103/PhysRevB.78.180502](https://doi.org/10.1103/PhysRevB.78.180502). URL: [https:](https://link.aps.org/doi/10.1103/PhysRevB.78.180502) [//link.aps.org/doi/10.1103/PhysRevB.78.180502](https://link.aps.org/doi/10.1103/PhysRevB.78.180502).
- <span id="page-15-5"></span>[6] Andreas Wallraff et al. "Strong coupling of a single photon to a superconducting qubit using circuit quantum electrodynamics". In: Nature 431 (Oct. 2004), pp. 162-7. doi: [10.1038/nature02851](https://doi.org/10.1038/nature02851).
- <span id="page-15-6"></span>[7] D. J. Egger and F. K. Wilhelm. "Multimode Circuit Quantum Electrodynamics with Hybrid Metamaterial Transmission Lines". In: Phys. Rev. Lett. 111 (16) Oct. 2013), p. 163601. DOI: [10.1103/PhysRevLett.111.163601](https://doi.org/10.1103/PhysRevLett.111.163601). URL: [https:](https://link.aps.org/doi/10.1103/PhysRevLett.111.163601) [//link.aps.org/doi/10.1103/PhysRevLett.111.163601](https://link.aps.org/doi/10.1103/PhysRevLett.111.163601).
- <span id="page-15-7"></span>[8] C. Caloz and T. Itoh. "Transmission line approach of left-handed (LH) materials and microstrip implementation of an artificial LH transmission line". In: IEEE Transactions on Antennas and Propagation  $52.5$  (2004), pp. 1159–1166. DOI: [10.1109/TAP.2004.827249](https://doi.org/10.1109/TAP.2004.827249).
- <span id="page-15-8"></span>[9] H. Wang et al. "Mode Structure in Superconducting Metamaterial Transmission-Line Resonators". In: *Phys. Rev. Applied* 11 (5 May 2019), p. 054062. DOI: [10.](https://doi.org/10.1103/PhysRevApplied.11.054062) [1103/PhysRevApplied.11.054062](https://doi.org/10.1103/PhysRevApplied.11.054062). url: [https://link.aps.org/doi/10.](https://link.aps.org/doi/10.1103/PhysRevApplied.11.054062) [1103/PhysRevApplied.11.054062](https://link.aps.org/doi/10.1103/PhysRevApplied.11.054062).
- <span id="page-15-9"></span>[10] Xiu Gu et al. "Microwave photonics with superconducting quantum circuits". In: *Physics Reports* 718-719 (July 2017), pp. 1–102. poi: [10.1016/j.physrep.](https://doi.org/10.1016/j.physrep.2017.10.002) [2017.10.002](https://doi.org/10.1016/j.physrep.2017.10.002).
- <span id="page-15-10"></span>[11] Morten Kjaergaard et al. "Superconducting Qubits: Current State of Play". In: Annual Review of Condensed Matter Physics 11.1 (Mar. 2020), pp. 369–395. ISSN: 1947-5462. DOI: 10.1146/annurev - conmatphys - 031119 - 050605. URL: <http://dx.doi.org/10.1146/annurev-conmatphys-031119-050605>.

<span id="page-16-0"></span>[12] P. Krantz et al. "A quantum engineer's guide to superconducting qubits". In: Applied Physics Reviews 6.2 (June 2019), p. 021318. ISSN: 1931-9401. DOI: [10.](https://doi.org/10.1063/1.5089550) [1063/1.5089550](https://doi.org/10.1063/1.5089550). url: <http://dx.doi.org/10.1063/1.5089550>.

# <span id="page-17-0"></span>A Appendix

Python design code

```
1 # -*-<i>coding</i>: <math>utf - 8</math> -*\frac{2}{2} """"
3 @author: Daniele Cucurachi
4
5 FULLY CUSTOMIZABLE COUPLED RESONATORS WAVEGUIDE GDS DESIGN
6
7 the function is divided into 4 parts, each part is related to the
      \rightarrow design of a specific component. In the end, all the
8 components are put together to crate the full WG. We have in order:
9
10 1) standard resonators pairs: the resonators pairs that form the
           \rightarrow wavequide
11
12 2) first resonator (+ feedline capacitor): first resonator in the
           \rightarrow waveguide and input feedline capacitor
13
14 3) last resonator: mirrored copy of the first resonator
15
16 4) ground design
17
18
19 NOTE: the various parts of the function create the negatives of the
        \rightarrow components designs we need. At the end, the boolean "not"
        \rightarrow function is used to subtract these negatives from a rectangle
          (which is the ground) and we obtain the full WG
        \rightarrow2021 \frac{1}{21} \frac{1}{21}22
23
24 #%% # -------------------------------------------------
25
26 # IMPORT LIBRARIES
27
28 import numpy as np
29 import gdspy
30
31
32 #%% # -------------------------------------------------
33
34
35
36 def RES_WG_APO():
37
```

```
38
39 # define all the parameters
40
41 # NOTE: the default dimensions are micrometers (um)
4243
44 """1) STANDARD RESONATORS PAIRS"""
45
46 # STD INDUCTOR
47
48 L = 279 # total length of the inductor wire
s = 4.5 # interspacing between the inductor windings
50 wid = 0.5 # width of inductor wire
51
52 compact = False # chose if compact inductor design or not
53
54 # STD CAPACITOR
55
56 A = 62 # horizontal dimension of U-capacitor
57 t = 3 # thickness of U-capacitor
58 add_t = 3 # added thickness for the bottom capacitor segment
59
60 # RESONATOR GROUND DISTANCES
61
62 dist_ground = 5 # distance between the capacitor and the ground
63 res_dist_1 = 29 # left lateral distance between two resonators in
       \rightarrow the resonators chain
64 res_dist_r = 29 # right lateral distance between two resonators in
       \rightarrow the resonators chain
65 inter_cell_dist = 4 # space between the two cells containing the
       \rightarrow resonators
66 ind_to_ground = 10 # last inductor seqment that connects the
       \rightarrow inductor to ground (inter_res_dist)
67 centre=(0,0) # centre where I start to draw the design
68
69
70 """2) FIRST RESONATOR + FEEDLINE CAPACITOR"""
71
72 # F INDUCTOR
73
74 Li = 279 # total length of the inductor wire
55 si = 4.5 # interspacing of turns of the inductor wire
76 wi = 0.5 # width of inductor wire
77 d_second =20 # last segment of the inductor (connection to ground.
       \rightarrow It defines the distance from ground at the top of the resonator
       \rightarrow cell)
```

```
78
```

```
79 compact ind = False # chose if compact inductor design or not
80 separated_elements = False # chose if you want to separate the
        \rightarrow inductor from the GND (for ANSYS simulations)
81 separation_length = 2 # separation length for separated elements
        \rightarrow option
82
83 # FEEDLINE INTERDIGITATED CAPACITOR
84
85 w2 = 2 # external feedline capacitor width
\text{sc} = 1.5 # lateral separation between the first resonator capacitor
       \rightarrow and the feedline capacitor
87 y = 2 # vertical separation between the first resonator capacitor
        \rightarrow and the feedline capacitor
88 w = 5 # feedline width
89 w1 = 4 # width of the internal capacitor arm
90 L2 = 45 # length of the internal capacitor arm
91 L1 = 40 # length of the external capacitor arm
92 Bwg = 200 # feedline length
93
94 # F CAPACITOR
95
96 Ac = 50 # horizontal dimension of U-capacitor (without additions,
       \rightarrow the final effective width is = A+Ws-t+Wl)
97 tc = 3 # thickness of U-capacitor (note that the thickness of the
       \rightarrow capacitor wide arm can be tuned individually)
98 Wl = t # wide arm thickness
99 Lw = 60 # wide arm length
100 Ws = (\text{sc} + \text{w2}) # thin arm shift
101
102 # GROUND DISTANCES
103
104 dist_ground_first = 65 # distance between capacitor and ground
105 first_dist_lateral_1 = 65 # left lateral distance between the
        \rightarrow capacitor and ground
106 first_dist_lateral_r = 28 # right lateral distance between the
        \rightarrow capacitor and ground
107 WG_width = 7 # feedline capacitor width
108
109 # define two parameters that will be used when creating the ground
        \rightarrow geometry
110
111 Lateral 1 = Bwg + w1 + sc112 Lateral_r = Ws + Ac + Wl - tc + first_dist_lateral_r
113
114
115
116 ""3) LAST RESONATOR"""
```

```
117
118
119
120 """A) GROUND DESIGN"""
121
122 Height = 250 # distance between the first cell and the edge of the
        \rightarrow ground piece
123 N_pairs = 3 # number of cells pairs in the waveguide (NOT including
        \rightarrow the first and last cells)
124
125
126
127 ""WANSYS/SONNET SIMULATION ADDED CAPACITOR THICKNESS """
128
129 # the parameter T1 offer the possibility to increase the thickness
        \rightarrow of the second and last second capacitors
130 # the parameter T2 offer the possibility to increase the thickness
        \rightarrow of the ground piece below the second and last second capacitors
131
132 # T1 and T2 are used to tune the capacitance to ground of the second
        \rightarrow and last second capacitors indipendently from the resonators in
        \rightarrow the waveguide
133
134 T1 = 0# thickness of the added capacitor piece
135
136 T2 = 0# thicnkess of the added ground piece
137
138
139
140
141 # -------------------------------------------------
142
143 # RESONATORS WAVEGUIDE DESIGN: PART 1) 2) 3) 4)
144
145
146
147
148 """1) STANDARD RESONATORS PAIRS"""
149
150
151 # CHECK FOR DIMENSIONS MISMATCH IN THE DISTANCES BETWEEN THE
        ,→ RESONATORS (STD RESONATORS CHAIN)
152
\frac{153}{153} while (res_dist_1 <= inter_cell_dist):
```

```
154 print("MISMATCH: res_dist_l (lateral distance between two
            \rightarrow resonators in the chain) is equal or smaller than
            \rightarrow <code>inter_cell_dist</code> (width of the ground in between two
            \rightarrow resonators cells)")
155 new = input("Enter new value for res_dist_l:")
156 res_dist_l = float(new)
157
158
\frac{159}{159} while (res_dist_r <= inter_cell_dist):
160 print("MISMATCH: res_dist_r (lateral distance between two
            \rightarrow resonators in the chain) is equal or smaller than
            \rightarrow inter_cell_dist (width of the ground in between two
            \rightarrow resonators)")
161 new = input("Enter new value for res_dist_r:")
res\_dist_r = float(new)163
164 dist_lateral_l = (res_dist_l - inter_cell_dist)/2
165 dist_lateral_r = (res_dist_r - inter_cell_dist)/2
166
167
168
169 # PARAMETERS CALCULATIONS
170
171 # horizontal dimension of the resonators pairs (these parameters
        \rightarrow will be used later to draw the wavequide geometry)
172
173 Horizonta = 2*A + 2*dist_1 \text{ateral}_1 + 2*dist_1 \text{ateral}_r +\rightarrow <code>inter_cell_dist</code>
174
175 half_res = A/2 + dist_lateral_l
176
177 #calculate horizontal dimension of inductor
178
179 b = 3/5*A - 2*t180
181 #calculate number of windings in the inductor
182
183 if compact == False:
184 N = int((L-(ind_to\_ground + (s+wid)))/(b+s))185 else:
N = round((L-(ind_to_ground + (s+wid)))/(b+s))187 \# \text{print('N: 'N)}188
189 #adapt first segment of the inductor
190
191 d_prime = (L - (N*(s+b)+ind_to\_ground)) + wid192
```

```
193 #calculate vertical dimension of capacitor
194
195 B = N*(s+wid) + d_prime + t
196
197 # define Ushape, set of points that will define the capacitor
       \rightarrow geometry
198
199 U=[]
200 x0, y0 = centre[0], centre[1]
201 U.append([x0, y0])
202
203 x1, y1 = x0, y0 - B + t/2
204 U.append([x1, y1])
205
206 x2, y2 = x1 + A - t, y1
207 U.append([x2, y2])
208
\begin{array}{rcl} \text{209} & \text{X3}, & \text{y3} = \text{x2}, & \text{y2} + \text{B} - \text{t/2} \end{array}210 U.append([x3, y3])
211
212 U = np.asarray(U)
213 centring = (A-t)/2, -B
214 U = U - centring
215
216 # additional thickness for the bottom segment of the capacitor
217
218 bottom_cap = gdspy.Rectangle((-A/2, -add_t), (+A/2,0))
219
220 #set of points that will define the meander (inductor) starting from
       \rightarrow centre of U
221
222 M = []
v0,w0 = centre[0], centre[1]+t
\frac{224}{224} M.append([v0,w0])
225 v1, ww1 = v0, w0 + d_prime - wid/2
226 M.append([v1,ww1])
227
228 for i in range(N):
229 if i\frac{9}{2} = 0:
230 v2, ww2 = M[-1][0] - b/2 + wid/2, M[-1][1]231 M.append([v2,ww2])
v3, w3 = v2, ww2 + (s+wid)233 M.append([v3,w3])
v4, w4 = v3 + (b/2 - wid/2), w3235 M.append([v4,w4])
\frac{236}{4} print(i)
2^{37} if (i+1)\%2 == 0:
```

```
238 v5, w5 = M[-1][0] + b/2 - wid/2, M[-1][1]
239 M.append([v5,w5])
v6, w6 = v5, w5 + (s+wid)241 M.append([v6,w6])
v7, w7 = v6 - b/2 + wid/2, w6
243 M.append([v7,w7])
\frac{1}{4} print(i)
245 if (i+1) == N:246 # print('end turns:', M[-1])
v8, w8 = M[-1][0], M[-1][1]v9, w9 = v8, w8 + ind_to_ground - wid/2249
250 # separate the inductor from the capacitor (OPTIONAL)
251
252 if separated_elements==True:
253 w9 = w9 - separation_length
254
255 M.append([v9,w9])
256
257
258
259 # GENERATE CAPACITOR AND INDUCTOR GEOMETRIES
260
261 Upath = gdspy. FlexPath(U,t) # capacitor
262 Mpath = gdspy. FlexPath(M, wid) # inductor
263
264 # create the negative that will be used later to create the final
       \rightarrow resonator wavequide
265
266 rec = gdspy.Rectangle((-(A/2) + dist_l = 1), -add_t
       \rightarrow -dist_ground), ((A/2) + dist_lateral_r, B + ind_to_ground -
       \rightarrow wid))
267
268 upneg = gdspy.boolean(rec, Upath, "not")
269 upneg = gdspy.boolean(upneg, Mpath, "not")
270 upneg = gdspy.boolean(upneg, bottom_cap, "not")
271
272 # create the negative of a pair of resonators spaced by
       \rightarrow inter_cell_dist
273
274 upneg_mirr = gdspy.copy(upneg) # copy the geometry
275 upneg_mirr.mirror([A/2 + dist_lateral_r + inter_cell_dist/2, 0],[A/2 + dist]\rightarrow + dist_lateral_r + inter_cell_dist/2, B]) # mirror the
       \rightarrow geometry
276
277
278
```

```
279
280 """2) FIRST RESONATOR"""
281
282
283 # CHECK IF THERE ARE DIMENSIONS MISMATCHES
284
285 while d_second < (y+w+(WG_width-w)/2):
286 print("ERROR: dimensions mismatch, d_second smaller than y+w")
287 new = input("Enter new value for d_second:")
288 d_second = float(new)289
290
291
292 # INDUCTOR DESIGN
293
294 #calculate horizontal dimension of inductor
295 # this inductor horizontal dimension will define the distance
       \rightarrow between the capacitors's arm and the inductor
296
297 bi = (3/5)*AC - 2*tc298
299 #calculate number of windings in the inductor
300
301 if compact_ind == False:
302 N = int((Li-(d_second + (si+wi)))/(bi+si))
303 else:
N = round((Li-(d\_second + (si+wi)))/(bi+si))305
306 #adapt start & end segment of inductor
307
308 d_prime = (Li - (N*(si+bi) + d_s) second) + wi
309
310 # define set of points that will be used to draw the inductor
311
312 M = []
313 v0, w0 = centre[0], centre[1]
314 M.append([v0,w0])
315 v1, ww1 = v0, w0 + d_prime - wi/2
316 M.append([v1,ww1])
317
318 # Inductor's windings
319
320 for i in range(N):
321
322 if i\frac{0}{2} = 0:
v2, wW2 = M[-1][0] + bi/2 - wi/2, M[-1][1]324 M.append([v2,ww2])
```

```
v3, w3 = v2, ww2 + (si+wi) # ww2 cause we already defined
             \rightarrow a w2 variable
326 M.append([v3,w3])
v4, w4 = v3 - (bi/2 - wi/2), w3328 M.append([v4,w4])
329
\int \sin^{-1} (i+1) \, \delta/2 = 0:
v5,w5 = M[-1][0] - bi/2 + wi/2, M[-1][1]332 M.append([v5,w5])
v6, w6 = v5, w5 + (si+wi)334 M.append([v6,w6])
v7, w7 = v6 + bi/2 - wi/2, w6336 M.append([v7,w7])
337
338 if (i+1) == N:
339 # print('end turns:', M[-1])
v8,w8 = M[-1][0], M[-1][1]v9, w9 = v8, w8 + d second + wi/2
342
343 # separate the inductor from the capacitor (OPTIONAL)
344
345 if separated_elements==True:
346 w9 = w9 - separation_length
347
348 M.append([v9,w9])
349
350 # shift the inductor (shift every point) and center it with respect
      \rightarrow to the capacitor (you drew the inductor centred in (0,0))
351
352 M = np.asarray(M)
353 centring = (Ac/2)+Ws, tc
354 M = M + centring
355
\frac{1}{356} # draw the inductor
357
358 ind = gdspy.FlexPath(M,wi) # FIRST RESONATOR INDUCTOR GEOMETRY
359
360
361
362
363 # CAPACITOR DESIGN
364
365 # Define length of the capacitor thin arm (the left one)
366
367 Lt = N*(si+wi) + d_prime + tc
368
369 # check for dimensions mismatch
```

```
370
371 while (Lw >= (Lt+d_second)):
372 print("MISMATCH: Lw is larger than Lt+d_second, it touches the
           \rightarrow ground")
373 new = input("Enter new value for Lw:")
374 Lw = float(new)
375
376 while ((L2-w+sc) > (y+(Lt-tc))):
377 print("ERROR: dimensions mismatch (L2 too large)")
378 new = input("Enter new value for L2:")
379 L2 = float(new)
380
381 # define the capacitor geometry
382
383 lcap = gdspy.Curve(0, 0).L(0, Lt, tc, Lt, tc, tc, Ws+Ac-tc, tc,,→ Ws+Ac-tc,tc, Ws+Ac-tc,Lw, Ws+Ac+Wl-tc,Lw, Ws+Ac+Wl-tc,0, 0,0)
384 cap = gdspy.Polygon(lcap.get_points()) # FIRST RESONATOR CAPACITOR
        ,→ GEOMETRY
385
386
387
388
389 """2) FEEDLINE CAPACITOR DESIGN"""
390
391
392 # DEFINE THE FEEDLINE CAPACITOR DESIGN
393
394 # the center (0,0) is the same used for the first resonator design
        \rightarrow function
395
396 wlcap = gdspy.Curve(-(sc+w1+Bwg),Lt+y).L(-(sc+w1+Bwg),Lt+y+w,
        \rightarrow tc+sc+w2,Lt+y+w, tc+sc+w2,Lt+y+w-L2, tc+sc,Lt+y+w-L2,
        \rightarrow tc+sc,Lt+y, -sc,Lt+y, -sc,Lt+y-L1, -(sc+w1),Lt+y-L1,
        \rightarrow -(sc+w1), Lt+y, -(sc+w1+Bwg), Lt+y)
397 wcap = gdspy.Polygon(wlcap.get_points())
398
399
400 # CREATE THE NEGATIVE of the whole resonator + waveguide capacitor
        \rightarrow geometry B
401
\frac{402}{402} contour_line = gdspy.Curve(-(sc+w1+first_dist_lateral_l),
```

```
403
           -dist_ground_first).L((Ws+Ac+Wl-tc+first_dist_lateral_r),(-dist_ground_first),
           (Ws+Ac+Wl-tc+first_dist_lateral_r),(Lt+d_second),
           -(sc+w1+first_dist_lateral_l),(Lt+d_second),
            -(sc+w1+first\_dist\_lateral\_l),(Lt+d\_second-(d\_second-(y+w+((WG\_width/2)- (w/2)))),
            -(sc+w1+Bwg), (Lt+d\_second-(d\_second-(y+w+((WG\_width/2) -(w/2)))),
           -(sc+w1+Bwg),(Lt+d_second-(d_second-(y+w+((WG_width/2) -
           (w/2))))-WG_width),
         \rightarrow -(sc+w1+first\_dist\_lateral\_l),(Lt+d\_second-(d\_second-(y+w+((WG\_width/2)- (w/2))))–WG_width),
         \rightarrow -(sc+wl+first\_dist\_lateral\_l),-(dist\_ground\_first))
         \hookrightarrow\rightarrow\rightarrow,→
        \rightarrow\hookrightarrow\hookrightarrow\hookrightarrow\hookrightarrow\rightarrow404 contour = gdspy.Polygon(contour_line.get_points())
405
406 # boolean subtarction to generate the negative
407
408 cont = gdspy.boolean(contour, wcap, "not")
409
_{410} co = gdspy.boolean(cont, cap, "not")
411
412 first_neg = gdspy.boolean(co, ind, "not") # NEGATIVE OF THE CELL
        ,→ CONTAINING FIRST RESONATOR + FEEDLINE CAPACITOR
413
414
415
416
417
418 """3) LAST RESONATOR: mirrored copy of the first one"""
419
420
421 # CREATE A MIRRORED COPY OF THE FIRST RESONATOR
422
423 first_neg_mirr = gdspy.copy(first_neg) # copy the geometry
424 first_neg_mirr.mirror([(2*Lateral_r + N_pairs*Horizonta +
         \rightarrow (N_pairs-1)*inter_cell_dist)/2, 0],[(2*Lateral_r +
         \rightarrow N_pairs*Horizonta + (N_pairs-1)*inter_cell_dist)/2, 1*1e-6]) #
         \rightarrow mirror the geometry
425
426 # NOTE: the mirrored copy is already shifted to the right at the end
        \rightarrow of the wavequide
427
428
429
430
\frac{431}{431} """ GROUND DESIGN"""
432
```

```
433
434 # CREATE THE GROUND RECTANGULAR GEOMETRY
435
436 # calculate the ground rectangle length
437
438 Rec_len = 2* \text{Lateral}_1 + 2* \text{Lateral}_r + N_-pairs* \text{Horizontal } +\rightarrow (N_pairs-1)*inter_cell_dist
439
440 #create the ground rectangle
441
442 ground = gdspy.Rectangle((-(Lateral_l),-(Height)), (Rec_len -
        \rightarrow Lateral_1, B + Height))
443
444
445 # SUBTRACT THE NEGATIVES
446
447 # create a cell array with the negatives of the standard resonators
        \rightarrow pairs
448
449 Res_pair_cell = gdspy.Cell("std_resonators_pair_reference")
450
451 Res_pair_cell.add(upneg)
452
453 Res_pair_cell.add(upneg_mirr)
454
455 res_array = gdspy.CellArray(ref_cell=Res_pair_cell, columns=N_pairs,
        \rightarrow rows=1, spacing = (Horizonta + inter_cell_dist, 0),
        \rightarrow origin=(Lateral_r + (half_res), 0)) # res_array type is cell
        \leftrightarrow array
456
457
458 # added capacitor thickness for ANSYS and Sonnet simulations
459 # check if there is a dimension mismatch
460
461 while ((T1+T2) > = dist\_ground):462 print("MISMATCH: ")
463 new = input("Enter new value for T1:")
_{464} T1 = float(new)
465 new = input("Enter new value for T2:")
_{466} T2 = float(new)
467
468
469 rec1 = gdspy.Rectangle((Lateral_r + dist_lateral_1,0),(Lateral_r +
        \rightarrow dist_lateral_1 + A,-T1))
470
```

```
r_4<sup>471</sup> rec2 = gdspy.Rectangle((Lateral_r +
        \rightarrow dist_lateral_1,-dist_ground),(Lateral_r + dist_lateral_1 + A,
           -dist\_ground + T2)) # this has this weird geometry cause I want
           to maximize the distance between the the first resonator
           capacitor and the ground
        \hookrightarrow\rightarrow\leftarrow472
473 rec3 = gdspy.Rectangle((Lateral_r + Horizonta*N_pairs +
        \rightarrow inter_cell_dist*(N_pairs-1) - (dist_lateral_l + A)
            ,0),(Lateral_r + Horizonta*N_pairs + inter_cell_dist*(N_pairs-1)
           - (dist_lateral_1 + A) + A, -T1))
        \hookrightarrow\hookrightarrow474
r = r + 475 rec4 = gdspy. Rectangle ((Lateral_r + Horizonta*N_pairs +
        \rightarrow inter_cell_dist*(N_pairs-1) - (dist_lateral_1 +
        \rightarrow A),-dist_ground), (Lateral_r + Horizonta*N_pairs +
        \rightarrow inter_cell_dist*(N_pairs-1) - (dist_lateral_1 + A) + A,
        \rightarrow -dist\_ground + T2))
476
477
478 res_array = gdspy.boolean(res_array, rec1, "not")
479 res_array = gdspy.boolean(res_array, rec2, "not")
480 res_array = gdspy.boolean(res_array, rec3, "not")
481 res_array = gdspy.boolean(res_array, rec4, "not")
482
483
484 # add the negatives of the components to one single cell (called
        \rightarrow "negatives")
485
486 negatives = gdspy.Cell("negatives")
487
488 negatives.add(res_array)
489
490 negatives.add(first_neg)
491
492 negatives.add(first_neg_mirr)
493
494
495 # subtract the negatives
496
497 WG = gdspy.boolean(ground, negatives, "not")
498
499 main = gdspy.Cell("Full_WG")
500 main.add(WG)
501
502
503 # RETURN: cell with the complete WG
504
505 return main
```

```
506
507
508 #%%# -------------------------------------------------
509
510 # PRINT A GDS FILE OF THE CHOSED DESIGN
511
512 main = RES_WG_APO()
513
514 lib = gdspy.GdsLibrary()
515
516 lib.add(main)
517
518 gdspy.LayoutViewer(lib)
519
520 # Save the library in a file called 'first.gds'
521
522 lib.write_gds('file_name.gds')
523
524
525
526
```